

Plane wave ground state expectation

Problem [1] 2.18 is, for a 1D SHO, show that

$$\langle 0 | e^{ikx} | 0 \rangle = \exp(-k^2 \langle 0 | x^2 | 0 \rangle / 2). \quad (1.1)$$

Despite the simple appearance of this problem, I found this quite involved to show. To do so, start with a series expansion of the expectation

$$\langle 0 | e^{ikx} | 0 \rangle = \sum_{m=0}^{\infty} \frac{(ik)^m}{m!} \langle 0 | x^m | 0 \rangle. \quad (1.2)$$

Let

$$X = (a + a^\dagger), \quad (1.3)$$

so that

$$x = \sqrt{\frac{\hbar}{2\omega m}} X = \frac{x_0}{\sqrt{2}} X. \quad (1.4)$$

Consider the first few values of $\langle 0 | X^n | 0 \rangle$

$$\begin{aligned} \langle 0 | X | 0 \rangle &= \langle 0 | (a + a^\dagger) | 0 \rangle \\ &= \langle 0 | 1 \rangle \\ &= 0, \end{aligned} \quad (1.5)$$

$$\begin{aligned} \langle 0 | X^2 | 0 \rangle &= \langle 0 | (a + a^\dagger)^2 | 0 \rangle \\ &= \langle 1 | 1 \rangle \\ &= 1, \end{aligned} \quad (1.6)$$

$$\begin{aligned} \langle 0 | X^3 | 0 \rangle &= \langle 0 | (a + a^\dagger)^3 | 0 \rangle \\ &= \langle 1 | (\sqrt{2} | 2 \rangle + | 0 \rangle) \rangle \\ &= 0. \end{aligned} \quad (1.7)$$

Whenever the power n in X^n is even, the bracket can be split into a bra that has only contributions from odd eigenstates and a ket with even eigenstates. We conclude that $\langle 0|X^n|0\rangle = 0$ when n is odd. Noting that $\langle 0|x^2|0\rangle = x_0^2/2$, this leaves

$$\begin{aligned}\langle 0|e^{ikx}|0\rangle &= \sum_{m=0}^{\infty} \frac{(ik)^{2m}}{(2m)!} \langle 0|x^{2m}|0\rangle \\ &= \sum_{m=0}^{\infty} \frac{(ik)^{2m}}{(2m)!} \left(\frac{x_0^2}{2}\right)^m \langle 0|X^{2m}|0\rangle \\ &= \sum_{m=0}^{\infty} \frac{1}{(2m)!} (-k^2 \langle 0|x^2|0\rangle)^m \langle 0|X^{2m}|0\rangle.\end{aligned}\tag{1.8}$$

This problem is now reduced to showing that

$$\frac{1}{(2m)!} \langle 0|X^{2m}|0\rangle = \frac{1}{m!2^m},\tag{1.9}$$

or

$$\begin{aligned}\langle 0|X^{2m}|0\rangle &= \frac{(2m)!}{m!2^m} \\ &= \frac{(2m)(2m-1)(2m-2)\cdots(2)(1)}{2^m m!} \\ &= \frac{2^m(m)(2m-1)(m-1)(2m-3)(m-2)\cdots(2)(3)(1)(1)}{2^m m!} \\ &= (2m-1)!!,\end{aligned}\tag{1.10}$$

where $n!! = n(n-2)(n-4)\cdots$.

It looks like $\langle 0|X^{2m}|0\rangle$ can be expanded by inserting an identity operator and proceeding recursively, like

$$\begin{aligned}\langle 0|X^{2m}|0\rangle &= \langle 0|X^2\left(\sum_{n=0}^{\infty}|n\rangle\langle n|\right)X^{2m-2}|0\rangle \\ &= \langle 0|X^2(|0\rangle\langle 0|+|2\rangle\langle 2|)X^{2m-2}|0\rangle \\ &= \langle 0|X^{2m-2}|0\rangle + \langle 0|X^2|2\rangle\langle 2|X^{2m-2}|0\rangle.\end{aligned}\tag{1.11}$$

This has made use of the observation that $\langle 0|X^2|n\rangle = 0$ for all $n \neq 0, 2$. The remaining term includes the factor

$$\begin{aligned}\langle 0|X^2|2\rangle &= \langle 0|(a+a^\dagger)^2|2\rangle \\ &= \left(\langle 0|+\sqrt{2}\langle 2|\right)|2\rangle \\ &= \sqrt{2},\end{aligned}\tag{1.12}$$

Since $\sqrt{2}|2\rangle = (a^\dagger)^2|0\rangle$, the expectation of interest can be written

$$\langle 0|X^{2m}|0\rangle = \langle 0|X^{2m-2}|0\rangle + \langle 0|a^2X^{2m-2}|0\rangle. \quad (1.13)$$

How do we expand the second term. Let's look at how a and X commute

$$\begin{aligned} aX &= [a, X] + Xa \\ &= [a, a + a^\dagger] + Xa \\ &= [a, a^\dagger] + Xa \\ &= 1 + Xa, \end{aligned} \quad (1.14)$$

$$\begin{aligned} a^2X &= a(aX) \\ &= a(1 + Xa) \\ &= a + aXa \\ &= a + (1 + Xa)a \\ &= 2a + Xa^2. \end{aligned} \quad (1.15)$$

Proceeding to expand a^2X^n we find

$$\begin{aligned} a^2X^3 &= 6X + 6X^2a + X^3a^2 \\ a^2X^4 &= 12X^2 + 8X^3a + X^4a^2 \\ a^2X^5 &= 20X^3 + 10X^4a + X^5a^2 \\ a^2X^6 &= 30X^4 + 12X^5a + X^6a^2. \end{aligned} \quad (1.16)$$

It appears that we have

$$[a^2X^n, X^na^2] = \beta_n X^{n-2} + 2nX^{n-1}a, \quad (1.17)$$

where

$$\beta_n = \beta_{n-1} + 2(n-1), \quad (1.18)$$

and $\beta_2 = 2$. Some goofing around shows that $\beta_n = n(n-1)$, so the induction hypothesis is

$$[a^2X^n, X^na^2] = n(n-1)X^{n-2} + 2nX^{n-1}a. \quad (1.19)$$

Let's check the induction

$$\begin{aligned} a^2X^{n+1} &= a^2X^nX \\ &= \left(n(n-1)X^{n-2} + 2nX^{n-1}a + X^na^2 \right) X \\ &= n(n-1)X^{n-1} + 2nX^{n-1}aX + X^na^2X \\ &= n(n-1)X^{n-1} + 2nX^{n-1}(1 + Xa) + X^n(2a + Xa^2) \\ &= n(n-1)X^{n-1} + 2nX^{n-1} + 2nX^na + 2X^na + X^{n+1}a^2 \\ &= X^{n+1}a^2 + (2 + 2n)X^na + (2n + n(n-1))X^{n-1} \\ &= X^{n+1}a^2 + 2(n+1)X^na + (n+1)nX^{n-1}, \end{aligned} \quad (1.20)$$

which concludes the induction, giving

$$\langle 0 | a^2 X^n | 0 \rangle = n(n-1) \langle 0 | X^{n-2} | 0 \rangle, \quad (1.21)$$

and

$$\langle 0 | X^{2m} | 0 \rangle = \langle 0 | X^{2m-2} | 0 \rangle + (2m-2)(2m-3) \langle 0 | X^{2m-4} | 0 \rangle. \quad (1.22)$$

Let

$$\sigma_n = \langle 0 | X^n | 0 \rangle, \quad (1.23)$$

so that the recurrence relation, for $2n \geq 4$ is

$$\sigma_{2n} = \sigma_{2n-2} + (2n-2)(2n-3)\sigma_{2n-4} \quad (1.24)$$

We want to show that this simplifies to

$$\sigma_{2n} = (2n-1)!! \quad (1.25)$$

The first values are

$$\sigma_0 = \langle 0 | X^0 | 0 \rangle = 1 \quad (1.26a)$$

$$\sigma_2 = \langle 0 | X^2 | 0 \rangle = 1 \quad (1.26b)$$

which gives us the right result for the first term in the induction

$$\begin{aligned} \sigma_4 &= \sigma_2 + 2 \times 1 \times \sigma_0 \\ &= 1 + 2 \\ &= 3!! \end{aligned} \quad (1.27)$$

For the general induction term, consider

$$\begin{aligned} \sigma_{2n+2} &= \sigma_{2n} + 2n(2n-1)\sigma_{2n-2} \\ &= (2n-1)!! + 2n(2n-1)(2n-3)!! \\ &= (2n+1)(2n-1)!! \\ &= (2n+1)!! \end{aligned} \quad (1.28)$$

which completes the final induction. That was also the last thing required to complete the proof, so we are done!

Bibliography

[1] Jun John Sakurai and Jim J Napolitano. *Modern quantum mechanics*. Pearson Higher Ed, 2014. 1