

Heisenberg picture spin precession

Exercise 1.1 Heisenberg picture spin precession ([1] pr. 2.1)

For the spin Hamiltonian

$$\begin{aligned} H &= -\frac{eB}{mc} S_z \\ &= \omega S_z, \end{aligned} \tag{1.1}$$

express and solve the Heisenberg equations of motion for $S_x(t)$, $S_y(t)$, and $S_z(t)$.

Answer for Exercise 1.1

The equations of motion are of the form

$$\begin{aligned} \frac{dS_i^H}{dt} &= \frac{1}{i\hbar} [S_i^H, H] \\ &= \frac{1}{i\hbar} [U^\dagger S_i U, H] \\ &= \frac{1}{i\hbar} (U^\dagger S_i U H - H U^\dagger S_i U) \\ &= \frac{1}{i\hbar} U^\dagger (S_i H - H S_i) U \\ &= \frac{\omega}{i\hbar} U^\dagger [S_i, S_z] U. \end{aligned} \tag{1.2}$$

These are

$$\begin{aligned} \frac{dS_x^H}{dt} &= -\omega U^\dagger S_y U \\ \frac{dS_y^H}{dt} &= \omega U^\dagger S_x U \\ \frac{dS_z^H}{dt} &= 0. \end{aligned} \tag{1.3}$$

To completely specify these equations, we need to expand $U(t)$, which is

$$\begin{aligned}
U(t) &= e^{-iHt/\hbar} \\
&= e^{-i\omega S_z t/\hbar} \\
&= e^{-i\omega\sigma_z t/2} \\
&= \cos(\omega t/2) - i\sigma_z \sin(\omega t/2) \\
&= \begin{bmatrix} \cos(\omega t/2) - i \sin(\omega t/2) & 0 \\ 0 & \cos(\omega t/2) + i \sin(\omega t/2) \end{bmatrix} \\
&= \begin{bmatrix} e^{-i\omega t/2} & 0 \\ 0 & e^{i\omega t/2} \end{bmatrix}.
\end{aligned} \tag{1.4}$$

The equations of motion can now be written out in full. To do so seems a bit silly since we also know that $S_x^H = U^\dagger S_x U$, $S_y^H = U^\dagger S_y U$. However, if that is temporarily forgotten, we can show that the Heisenberg equations of motion can be solved for these too.

$$\begin{aligned}
U^\dagger S_x U &= \frac{\hbar}{2} \begin{bmatrix} e^{i\omega t/2} & 0 \\ 0 & e^{-i\omega t/2} \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} e^{-i\omega t/2} & 0 \\ 0 & e^{i\omega t/2} \end{bmatrix} \\
&= \frac{\hbar}{2} \begin{bmatrix} 0 & e^{i\omega t/2} \\ e^{-i\omega t/2} & 0 \end{bmatrix} \begin{bmatrix} e^{-i\omega t/2} & 0 \\ 0 & e^{i\omega t/2} \end{bmatrix} \\
&= \frac{\hbar}{2} \begin{bmatrix} 0 & e^{i\omega t} \\ e^{-i\omega t} & 0 \end{bmatrix},
\end{aligned} \tag{1.5}$$

and

$$\begin{aligned}
U^\dagger S_y U &= \frac{\hbar}{2} \begin{bmatrix} e^{i\omega t/2} & 0 \\ 0 & e^{-i\omega t/2} \end{bmatrix} \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} e^{-i\omega t/2} & 0 \\ 0 & e^{i\omega t/2} \end{bmatrix} \\
&= \frac{i\hbar}{2} \begin{bmatrix} 0 & -e^{i\omega t/2} \\ e^{-i\omega t/2} & 0 \end{bmatrix} \begin{bmatrix} e^{-i\omega t/2} & 0 \\ 0 & e^{i\omega t/2} \end{bmatrix} \\
&= \frac{i\hbar}{2} \begin{bmatrix} 0 & -e^{i\omega t} \\ e^{-i\omega t} & 0 \end{bmatrix}.
\end{aligned} \tag{1.6}$$

The equations of motion are now fully specified

$$\begin{aligned}
\frac{dS_x^H}{dt} &= -\frac{i\hbar\omega}{2} \begin{bmatrix} 0 & -e^{i\omega t} \\ e^{-i\omega t} & 0 \end{bmatrix} \\
\frac{dS_y^H}{dt} &= \frac{\hbar\omega}{2} \begin{bmatrix} 0 & e^{i\omega t} \\ e^{-i\omega t} & 0 \end{bmatrix} \\
\frac{dS_z^H}{dt} &= 0.
\end{aligned} \tag{1.7}$$

Integration gives

$$\begin{aligned}
S_x^H &= \frac{\hbar}{2} \begin{bmatrix} 0 & e^{i\omega t} \\ e^{-i\omega t} & 0 \end{bmatrix} + C \\
S_y^H &= \frac{\hbar}{2} \begin{bmatrix} 0 & -ie^{i\omega t} \\ ie^{-i\omega t} & 0 \end{bmatrix} + C \\
S_z^H &= C.
\end{aligned} \tag{1.8}$$

The integration constants are fixed by the boundary condition $S_i^H(0) = S_i$, so

$$\begin{aligned}
S_x^H &= \frac{\hbar}{2} \begin{bmatrix} 0 & e^{i\omega t} \\ e^{-i\omega t} & 0 \end{bmatrix} \\
S_y^H &= \frac{i\hbar}{2} \begin{bmatrix} 0 & -e^{i\omega t} \\ e^{-i\omega t} & 0 \end{bmatrix} \\
S_z^H &= S_z.
\end{aligned} \tag{1.9}$$

Observe that these integrated values S_x^H, S_y^H match eq. (1.5), and eq. (1.6) as expected.

Bibliography

- [1] Jun John Sakurai and Jim J Napolitano. *Modern quantum mechanics*. Pearson Higher Ed, 2014. 1.1