

Time evolution of spin half probability and dispersion

Exercise 1.1 Time evolution of spin half probability and dispersion ([1] pr. 2.3)

A spin 1/2 system $\mathbf{S} \cdot \hat{\mathbf{n}}$, with $\hat{\mathbf{n}} = \sin \beta \hat{\mathbf{x}} + \cos \beta \hat{\mathbf{z}}$, is in state with eigenvalue $\hbar/2$, acted on by a magnetic field of strength B in the $+z$ direction.

1. If S_x is measured at time t , what is the probability of getting $+\hbar/2$?
2. Evaluate the dispersion in S_x as a function of t , that is,

$$\langle (S_x - \langle S_x \rangle)^2 \rangle. \quad (1.1)$$

3. Check your answers for $\beta \rightarrow 0, \pi/2$ to see if they make sense.

Answer for Exercise 1.1

Part 1. The spin operator in matrix form is

$$\begin{aligned} \mathbf{S} \cdot \hat{\mathbf{n}} &= \frac{\hbar}{2} (\sigma_z \cos \beta + \sigma_x \sin \beta) \\ &= \frac{\hbar}{2} \left(\begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \cos \beta + \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \sin \beta \right) \\ &= \frac{\hbar}{2} \begin{bmatrix} \cos \beta & \sin \beta \\ \sin \beta & -\cos \beta \end{bmatrix}. \end{aligned} \quad (1.2)$$

The $|\mathbf{S} \cdot \hat{\mathbf{n}}; +\rangle$ eigenstate is found from

$$(\mathbf{S} \cdot \hat{\mathbf{n}} - \hbar/2) \begin{bmatrix} a \\ b \end{bmatrix} = 0, \quad (1.3)$$

or

$$\begin{aligned} 0 &= (\cos \beta - 1) a + \sin \beta b \\ &= (-2 \sin^2(\beta/2)) a + 2 \sin(\beta/2) \cos(\beta/2) b \\ &= (-\sin(\beta/2)) a + \cos(\beta/2) b, \end{aligned} \quad (1.4)$$

or

$$|S \cdot \hat{\mathbf{n}}; +\rangle = \begin{bmatrix} \cos(\beta/2) \\ \sin(\beta/2) \end{bmatrix}. \quad (1.5)$$

The Hamiltonian is

$$H = -\frac{eB}{mc} S_z = -\frac{eB\hbar}{2mc} \sigma_z, \quad (1.6)$$

so the time evolution operator is

$$U = e^{-iHt/\hbar} = e^{\frac{ieBt}{2mc} \sigma_z}. \quad (1.7)$$

Let $\omega = eB/(2mc)$, so

$$\begin{aligned} U &= e^{i\sigma_z \omega t} \\ &= \cos(\omega t) + i\sigma_z \sin(\omega t) \\ &= \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \cos(\omega t) + i \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \sin(\omega t) \\ &= \begin{bmatrix} e^{i\omega t} & 0 \\ 0 & e^{-i\omega t} \end{bmatrix}. \end{aligned} \quad (1.8)$$

The time evolution of the initial state is

$$\begin{aligned} |S \cdot \hat{\mathbf{n}}; +\rangle (t) &= U |S \cdot \hat{\mathbf{n}}; +\rangle (0) \\ &= \begin{bmatrix} e^{i\omega t} & 0 \\ 0 & e^{-i\omega t} \end{bmatrix} \begin{bmatrix} \cos(\beta/2) \\ \sin(\beta/2) \end{bmatrix} \\ &= \begin{bmatrix} \cos(\beta/2)e^{i\omega t} \\ \sin(\beta/2)e^{-i\omega t} \end{bmatrix}. \end{aligned} \quad (1.9)$$

The probability of finding the state in $|S \cdot \hat{\mathbf{x}}; +\rangle$ at time t (i.e. measuring S_x and finding $\hbar/2$) is

$$\begin{aligned} |\langle S \cdot \hat{\mathbf{x}}; + | S \cdot \hat{\mathbf{n}}; + \rangle|^2 &= \left| \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \end{bmatrix} \begin{bmatrix} \cos(\beta/2)e^{i\omega t} \\ \sin(\beta/2)e^{-i\omega t} \end{bmatrix} \right|^2 \\ &= \frac{1}{2} \left| \cos(\beta/2)e^{i\omega t} + \sin(\beta/2)e^{-i\omega t} \right|^2 \\ &= \frac{1}{2} (1 + 2 \cos(\beta/2) \sin(\beta/2) \cos(2\omega t)) \\ &= \frac{1}{2} (1 + \sin(\beta) \cos(2\omega t)). \end{aligned} \quad (1.10)$$

Part 2. To calculate the dispersion first note that

$$\begin{aligned} S_x^2 &= \left(\frac{\hbar}{2}\right)^2 \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}^2 \\ &= \left(\frac{\hbar}{2}\right)^2 \cdot \end{aligned} \quad (1.11)$$

so only the first order expectation is non-trivial to calculate. That is

$$\begin{aligned} \langle S_x \rangle &= \frac{\hbar}{2} \begin{bmatrix} \cos(\beta/2)e^{-i\omega t} & \sin(\beta/2)e^{i\omega t} \end{bmatrix} \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} \cos(\beta/2)e^{i\omega t} \\ \sin(\beta/2)e^{-i\omega t} \end{bmatrix} \\ &= \frac{\hbar}{2} \begin{bmatrix} \cos(\beta/2)e^{-i\omega t} & \sin(\beta/2)e^{i\omega t} \end{bmatrix} \begin{bmatrix} \sin(\beta/2)e^{-i\omega t} \\ \cos(\beta/2)e^{i\omega t} \end{bmatrix} \\ &= \frac{\hbar}{2} \sin(\beta/2) \cos(\beta/2) (e^{-2i\omega t} + e^{2i\omega t}) \\ &= \frac{\hbar}{2} \sin \beta \cos(2\omega t). \end{aligned} \quad (1.12)$$

This gives

$$\langle (\Delta S_x)^2 \rangle = \left(\frac{\hbar}{2}\right)^2 (1 - \sin^2 \beta \cos^2(2\omega t)) \quad (1.13)$$

Part 3. For $\beta = 0$, $\hat{n} = \hat{z}$, and $\beta = \pi/2$, $\hat{n} = \hat{x}$. For the first case, the state is in an eigenstate of S_z , so must evolve as

$$|S \cdot \hat{n}; +\rangle (t) = |S \cdot \hat{n}; +\rangle (0)e^{i\omega t}. \quad (1.14)$$

The probability of finding it in state $|S \cdot \hat{x}; +\rangle$ is therefore

$$\begin{aligned} &\left| \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \end{bmatrix} \begin{bmatrix} e^{i\omega t} \\ 0 \end{bmatrix} \right|^2 \\ &= \frac{1}{2} |e^{i\omega t}|^2 \\ &= \frac{1}{2} \\ &= \frac{1}{2} (1 + \sin(0) \cos(2\omega t)). \end{aligned} \quad (1.15)$$

This matches eq. (1.10) as expected.

For $\beta = \pi/2$ we have

$$\begin{aligned} |S \cdot \hat{x}; +\rangle (t) &= \frac{1}{\sqrt{2}} \begin{bmatrix} e^{i\omega t} & 0 \\ 0 & e^{-i\omega t} \end{bmatrix} \begin{bmatrix} 1 \\ 1 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} \begin{bmatrix} e^{i\omega t} \\ e^{-i\omega t} \end{bmatrix}. \end{aligned} \quad (1.16)$$

The probability for the $\hbar/2 S_x$ measurement at time t is

$$\begin{aligned} \left| \frac{1}{2} \begin{bmatrix} 1 & 1 \\ e^{i\omega t} & e^{-i\omega t} \end{bmatrix} \right|^2 &= \frac{1}{4} |e^{i\omega t} + e^{-i\omega t}|^2 \\ &= \cos^2(\omega t) \\ &= \frac{1}{2} (1 + \sin(\pi/2) \cos(2\omega t)). \end{aligned} \tag{1.17}$$

Again, this matches the expected value.

For the dispersions, at $\beta = 0$, the dispersion is

$$\left(\frac{\hbar}{2} \right)^2 \tag{1.18}$$

This is the maximum dispersion, which makes sense since we are measuring S_x when the initial state is $|S \cdot \hat{z}; +\rangle$. For $\beta = \pi/2$ the dispersion is

$$\left(\frac{\hbar}{2} \right)^2 \sin^2(2\omega t). \tag{1.19}$$

This starts off as zero dispersion (because the initial state is $|S \cdot \hat{x}; +\rangle$), but then oscillates.

Bibliography

- [1] Jun John Sakurai and Jim J Napolitano. *Modern quantum mechanics*. Pearson Higher Ed, 2014. 1.1