

## Commutators for some symmetry operators

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*Q: [1] pr 4.2* If  $\mathcal{T}_{\mathbf{d}}$ ,  $\mathcal{D}(\hat{\mathbf{n}}, \phi)$ , and  $\pi$  denote the translation, rotation, and parity operators respectively. Which of the following commute and why

- (a)  $\mathcal{T}_{\mathbf{d}}$  and  $\mathcal{T}_{\mathbf{d}'}$ , translations in different directions.
- (b)  $\mathcal{D}(\hat{\mathbf{n}}, \phi)$  and  $\mathcal{D}(\hat{\mathbf{n}}', \phi')$ , rotations in different directions.
- (c)  $\mathcal{T}_{\mathbf{d}}$  and  $\pi$ .
- (d)  $\mathcal{D}(\hat{\mathbf{n}}, \phi)$  and  $\pi$ .

*A: (a)* Consider

$$\begin{aligned}\mathcal{T}_{\mathbf{d}}\mathcal{T}_{\mathbf{d}'}|\mathbf{x}\rangle &= \mathcal{T}_{\mathbf{d}}|\mathbf{x} + \mathbf{d}'\rangle \\ &= |\mathbf{x} + \mathbf{d}' + \mathbf{d}\rangle,\end{aligned}\tag{1.1}$$

and the reverse application of the translation operators

$$\begin{aligned}\mathcal{T}_{\mathbf{d}'}\mathcal{T}_{\mathbf{d}}|\mathbf{x}\rangle &= \mathcal{T}_{\mathbf{d}'}|\mathbf{x} + \mathbf{d}\rangle \\ &= |\mathbf{x} + \mathbf{d} + \mathbf{d}'\rangle \\ &= |\mathbf{x} + \mathbf{d}' + \mathbf{d}\rangle.\end{aligned}\tag{1.2}$$

so we see that

$$[\mathcal{T}_{\mathbf{d}}, \mathcal{T}_{\mathbf{d}'}]|\mathbf{x}\rangle = 0,\tag{1.3}$$

for any position state  $|\mathbf{x}\rangle$ , and therefore in general they commute.

*A: (b)* That rotations do not commute when they are in different directions (like any two orthogonal directions) need not be belaboured.

A: (c) We have

$$\begin{aligned}\mathcal{T}_{\mathbf{d}}\pi|\mathbf{x}\rangle &= \mathcal{T}_{\mathbf{d}}|-\mathbf{x}\rangle \\ &= |-\mathbf{x} + \mathbf{d}\rangle,\end{aligned}\tag{1.4}$$

yet

$$\begin{aligned}\pi\mathcal{T}_{\mathbf{d}}|\mathbf{x}\rangle &= \pi|\mathbf{x} + \mathbf{d}\rangle \\ &= |-\mathbf{x} - \mathbf{d}\rangle \\ &\neq |-\mathbf{x} + \mathbf{d}\rangle.\end{aligned}\tag{1.5}$$

so, in general  $[\mathcal{T}_{\mathbf{d}}, \pi] \neq 0$ .

A: (d) We have

$$\begin{aligned}\pi\mathcal{D}(\hat{\mathbf{n}}, \phi)|\mathbf{x}\rangle &= \pi\mathcal{D}(\hat{\mathbf{n}}, \phi)\pi^\dagger\pi|\mathbf{x}\rangle \\ &= \pi\mathcal{D}(\hat{\mathbf{n}}, \phi)\pi^\dagger\pi|\mathbf{x}\rangle \\ &= \pi\left(\sum_{k=0}^{\infty}\frac{(-i\mathbf{J}\cdot\hat{\mathbf{n}})^k}{k!}\right)\pi^\dagger\pi|\mathbf{x}\rangle \\ &= \sum_{k=0}^{\infty}\frac{(-i(\pi\mathbf{J}\pi^\dagger)\cdot(\pi\hat{\mathbf{n}}\pi^\dagger))^k}{k!}\pi|\mathbf{x}\rangle \\ &= \sum_{k=0}^{\infty}\frac{(-i\mathbf{J}\cdot\hat{\mathbf{n}})^k}{k!}\pi|\mathbf{x}\rangle \\ &= \mathcal{D}(\hat{\mathbf{n}}, \phi)\pi|\mathbf{x}\rangle,\end{aligned}\tag{1.6}$$

so  $[\mathcal{D}(\hat{\mathbf{n}}, \phi), \pi]|\mathbf{x}\rangle = 0$ , for any position state  $|\mathbf{x}\rangle$ , and therefore these operators commute in general.

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## Bibliography

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- [1] Jun John Sakurai and Jim J Napolitano. *Modern quantum mechanics*. Pearson Higher Ed, 2014. 1