

Plane wave and spinor under time reversal

Q: [1] pr 4.7

- (a) Find the time reversed form of a spinless plane wave state in three dimensions.
- (b) For the eigenspinor of $\boldsymbol{\sigma} \cdot \hat{\mathbf{n}}$ expressed in terms of polar and azimuthal angles β and γ , show that $-i\sigma_y\chi^*(\hat{\mathbf{n}})$ has the reversed spin direction.

A: part (a) The Hamiltonian for a plane wave is

$$H = \frac{\mathbf{p}^2}{2m} = i\frac{\partial}{\partial t} \quad (1.1)$$

Under time reversal the momentum side transforms as

$$\begin{aligned} \Theta \frac{\mathbf{p}^2}{2m} \Theta^{-1} &= \frac{(\Theta \mathbf{p} \Theta^{-1}) \cdot (\Theta \mathbf{p} \Theta^{-1})}{2m} \\ &= \frac{(-\mathbf{p}) \cdot (-\mathbf{p})}{2m} \\ &= \frac{\mathbf{p}^2}{2m}. \end{aligned} \quad (1.2)$$

The time derivative side of the equation is also time reversal invariant

$$\begin{aligned} \Theta i \frac{\partial}{\partial t} \Theta^{-1} &= \Theta i \Theta^{-1} \Theta \frac{\partial}{\partial t} \Theta^{-1} \\ &= -i \frac{\partial}{\partial(-t)} \\ &= i \frac{\partial}{\partial t}. \end{aligned} \quad (1.3)$$

Solutions to this equation are linear combinations of

$$\psi(\mathbf{x}, t) = e^{i\mathbf{k} \cdot \mathbf{x} - iEt/\hbar}, \quad (1.4)$$

where $\hbar^2 \mathbf{k}^2 / 2m = E$, the energy of the particle. Under time reversal we have

$$\begin{aligned}
\psi(\mathbf{x}, t) &\rightarrow e^{-i\mathbf{k}\cdot\mathbf{x}+iE(-t)/\hbar} \\
&= \left(e^{i\mathbf{k}\cdot\mathbf{x}-iE(-t)/\hbar} \right)^* \\
&= \psi^*(\mathbf{x}, -t)
\end{aligned} \tag{1.5}$$

A: part (b) The text uses a requirement for time reversal of spin states to show that the Pauli matrix form of the time reversal operator is

$$\Theta = -i\sigma_y K, \tag{1.6}$$

where K is a complex conjugating operator. The form of the spin up state used in that demonstration was

$$\begin{aligned}
|\hat{\mathbf{n}}; +\rangle &= e^{-iS_z\beta/\hbar} e^{-iS_y\gamma/\hbar} |+\rangle \\
&= e^{-i\sigma_z\beta/2} e^{-i\sigma_y\gamma/2} |+\rangle \\
&= (\cos(\beta/2) - i\sigma_z \sin(\beta/2)) (\cos(\gamma/2) - i\sigma_y \sin(\gamma/2)) |+\rangle \\
&= \left(\cos(\beta/2) - i \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \sin(\beta/2) \right) \left(\cos(\gamma/2) - i \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \sin(\gamma/2) \right) |+\rangle \\
&= \begin{bmatrix} e^{-i\beta/2} & 0 \\ 0 & e^{i\beta/2} \end{bmatrix} \begin{bmatrix} \cos(\gamma/2) & -\sin(\gamma/2) \\ \sin(\gamma/2) & \cos(\gamma/2) \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \\
&= \begin{bmatrix} e^{-i\beta/2} & 0 \\ 0 & e^{i\beta/2} \end{bmatrix} \begin{bmatrix} \cos(\gamma/2) \\ \sin(\gamma/2) \end{bmatrix} \\
&= \begin{bmatrix} \cos(\gamma/2)e^{-i\beta/2} \\ \sin(\gamma/2)e^{i\beta/2} \end{bmatrix}.
\end{aligned} \tag{1.7}$$

The state orthogonal to this one is claimed to be

$$\begin{aligned}
|\hat{\mathbf{n}}; -\rangle &= e^{-iS_z\beta/\hbar} e^{-iS_y(\gamma+\pi)/\hbar} |+\rangle \\
&= e^{-i\sigma_z\beta/2} e^{-i\sigma_y(\gamma+\pi)/2} |+\rangle.
\end{aligned} \tag{1.8}$$

We have

$$\cos((\gamma + \pi)/2) = \operatorname{Re} e^{i(\gamma+\pi)/2} = \operatorname{Re} i e^{i\gamma/2} = -\sin(\gamma/2), \tag{1.9}$$

and

$$\sin((\gamma + \pi)/2) = \operatorname{Im} e^{i(\gamma+\pi)/2} = \operatorname{Im} i e^{i\gamma/2} = \cos(\gamma/2), \tag{1.10}$$

so we should have

$$|\hat{\mathbf{n}}; -\rangle = \begin{bmatrix} -\sin(\gamma/2)e^{-i\beta/2} \\ \cos(\gamma/2)e^{i\beta/2} \end{bmatrix}. \tag{1.11}$$

This looks right, but we can sanity check orthogonality

$$\langle \hat{\mathbf{n}}; - | \hat{\mathbf{n}}; + \rangle = \begin{bmatrix} -\sin(\gamma/2)e^{i\beta/2} & \cos(\gamma/2)e^{-i\beta/2} \end{bmatrix} \begin{bmatrix} \cos(\gamma/2)e^{-i\beta/2} \\ \sin(\gamma/2)e^{i\beta/2} \end{bmatrix} = 0, \quad (1.12)$$

as expected.

The task at hand appears to be the operation on the column representation of $|\hat{\mathbf{n}}; +\rangle$ using the Pauli representation of the time reversal operator. That is

$$\begin{aligned} \Theta |\hat{\mathbf{n}}; +\rangle &= -i\sigma_y K \begin{bmatrix} e^{-i\beta/2} \cos(\gamma/2) \\ e^{i\beta/2} \sin(\gamma/2) \end{bmatrix} \\ &= -i \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \begin{bmatrix} e^{i\beta/2} \cos(\gamma/2) \\ e^{-i\beta/2} \sin(\gamma/2) \end{bmatrix} \\ &= \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} \begin{bmatrix} e^{i\beta/2} \cos(\gamma/2) \\ e^{-i\beta/2} \sin(\gamma/2) \end{bmatrix} \\ &= \begin{bmatrix} -e^{-i\beta/2} \sin(\gamma/2) \\ e^{i\beta/2} \cos(\gamma/2) \end{bmatrix} \\ &= |\hat{\mathbf{n}}; -\rangle. \quad \square \end{aligned} \quad (1.13)$$

Bibliography

[1] Jun John Sakurai and Jim J Napolitano. *Modern quantum mechanics*. Pearson Higher Ed, 2014. 1