Spherical gradient, divergence, curl and Laplacian

Unit vectors Two of the spherical unit vectors we can immediately write by inspection.

$$\hat{\mathbf{r}} = \mathbf{e}_1 \sin \theta \cos \phi + \mathbf{e}_2 \sin \theta \sin \phi + \mathbf{e}_3 \cos \theta$$

$$\hat{\boldsymbol{\phi}} = -\mathbf{e}_1 \sin \theta + \mathbf{e}_2 \cos \phi$$
(1.1)

We can compute $\hat{\theta}$ by utilizing the right hand triplet property

$$\hat{\boldsymbol{\theta}} = \hat{\boldsymbol{\phi}} \times \hat{\mathbf{r}}
= \begin{vmatrix} \mathbf{e}_{1} & \mathbf{e}_{2} & \mathbf{e}_{3} \\ -S_{\phi} & C_{\phi} & 0 \\ S_{\theta}C_{\phi} & S_{\theta}S_{\phi} & C_{\theta} \end{vmatrix}
= \mathbf{e}_{1} \left(C_{\theta}C_{\phi} \right) + \mathbf{e}_{2} \left(C_{\theta}S_{\phi} \right) + \mathbf{e}_{3} \left(-S_{\theta} \left(S_{\phi}^{2} + C_{\phi}^{2} \right) \right)
= \mathbf{e}_{1} \cos \theta \cos \phi + \mathbf{e}_{2} \cos \theta \sin \phi - \mathbf{e}_{3} \sin \theta.$$
(1.2)

Here I've used $C_{\theta} = \cos \theta$, $S_{\phi} = \sin \phi$, \cdots as a convenient shorthand. Observe that with $i = \mathbf{e}_1 \mathbf{e}_2$, these unit vectors admit a small factorization that makes further manipulation easier

$$\hat{\mathbf{r}} = \mathbf{e}_1 e^{i\phi} \sin \theta + \mathbf{e}_3 \cos \theta$$

$$\hat{\boldsymbol{\theta}} = \cos \theta \mathbf{e}_1 e^{i\phi} - \sin \theta \mathbf{e}_3$$

$$\hat{\boldsymbol{\phi}} = \mathbf{e}_2 e^{i\phi}$$
(1.3)

It should also be the case that $\hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\hat{\boldsymbol{\phi}}=I$, where $I=\mathbf{e}_1\mathbf{e}_2\mathbf{e}_3=\mathbf{e}_{123}$ is the \mathbb{R}^3 pseudoscalar, which is straightforward to check

$$\hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\hat{\boldsymbol{\phi}} = \left(\mathbf{e}_{1}e^{i\phi}\sin\theta + \mathbf{e}_{3}\cos\theta\right)\left(\cos\theta\mathbf{e}_{1}e^{i\phi} - \sin\theta\mathbf{e}_{3}\right)\mathbf{e}_{2}e^{i\phi}
= \left(\sin\theta\cos\theta - \cos\theta\sin\theta + \mathbf{e}_{31}e^{i\phi}\left(\cos^{2}\theta + \sin^{2}\theta\right)\right)\mathbf{e}_{2}e^{i\phi}
= \mathbf{e}_{31}\mathbf{e}_{2}e^{-i\phi}e^{i\phi}
= \mathbf{e}_{123}.$$
(1.4)

This property could also have been used to compute $\hat{\theta}$.

Gradient To compute the gradient, note that the coordinate vectors for the spherical parameterization are

$$\mathbf{x}_{r} = \frac{\partial \mathbf{r}}{\partial r}$$

$$= \frac{\partial (r\hat{\mathbf{r}})}{\partial r}$$

$$= \hat{\mathbf{r}} + r \frac{\partial \hat{\mathbf{r}}}{\partial r}$$

$$= \hat{\mathbf{r}}, \qquad (1.5a)$$

$$\mathbf{x}_{\theta} = \frac{\partial (r\hat{\mathbf{r}})}{\partial \theta}$$

$$= r \frac{\partial}{\partial \theta} \left(S_{\theta} \mathbf{e}_{1} e^{i\phi} + C_{\theta} \mathbf{e}_{3} \right)$$

$$= r \frac{\partial}{\partial \theta} \left(C_{\theta} \mathbf{e}_{1} e^{i\phi} - S_{\theta} \mathbf{e}_{3} \right)$$

$$= r \hat{\boldsymbol{\theta}},$$
(1.5b)

$$\mathbf{x}_{\phi} = \frac{\partial(r\hat{\mathbf{r}})}{\partial\phi}$$

$$= r\frac{\partial}{\partial\phi} \left(S_{\theta} \mathbf{e}_{1} e^{i\phi} + C_{\theta} \mathbf{e}_{3} \right)$$

$$= rS_{\theta} \mathbf{e}_{2} e^{i\phi}$$

$$= r\sin\theta \hat{\boldsymbol{\phi}}.$$
(1.5c)

Since these are all normal, the dual vectors defined by $\mathbf{x}^j \cdot \mathbf{x}_k = \delta_k^j$, can be obtained by inspection

$$\mathbf{x}^{r} = \hat{\mathbf{r}}$$

$$\mathbf{x}^{\theta} = \frac{1}{r}\hat{\boldsymbol{\theta}}$$

$$\mathbf{x}^{\phi} = \frac{1}{r\sin\theta}\hat{\boldsymbol{\phi}}.$$
(1.6)

The gradient follows immediately

$$\nabla = \mathbf{x}^r \frac{\partial}{\partial r} + \mathbf{x}^\theta \frac{\partial}{\partial \theta} + \mathbf{x}^\phi \frac{\partial}{\partial \hat{\boldsymbol{\phi}}},$$
(1.7)

or

$$\nabla = \hat{\mathbf{r}} \frac{\partial}{\partial r} + \frac{\hat{\boldsymbol{\theta}}}{r} \frac{\partial}{\partial \theta} + \frac{\hat{\boldsymbol{\phi}}}{r \sin \theta} \frac{\partial}{\partial \hat{\boldsymbol{\phi}}}.$$
 (1.8)

More information on this general dual-vector technique of computing the gradient in curvilinear coordinate systems can be found in [2].

Partials To compute the divergence, curl and Laplacian, we'll need the partials of each of the unit vectors $\partial \hat{\mathbf{r}}/\partial \theta$, $\partial \hat{\mathbf{r}}/\partial \phi$, $\partial \hat{\boldsymbol{\theta}}/\partial \theta$, $\partial \hat{\boldsymbol{\theta}}/\partial \phi$, $\partial \hat{\boldsymbol{\theta}}/\partial \phi$.

The $\hat{\theta}$ partials are

$$\frac{\partial \hat{\boldsymbol{\theta}}}{\partial \theta} = \frac{\partial}{\partial \theta} \left(C_{\theta} \mathbf{e}_{1} e^{i\phi} - S_{\theta} \mathbf{e}_{3} \right)
= -S_{\theta} \mathbf{e}_{1} e^{i\phi} - C_{\theta} \mathbf{e}_{3}
= -\hat{\mathbf{r}},$$
(1.9)

$$\frac{\partial \hat{\boldsymbol{\theta}}}{\partial \phi} = \frac{\partial}{\partial \phi} \left(C_{\theta} \mathbf{e}_{1} e^{i\phi} - S_{\theta} \mathbf{e}_{3} \right)
= C_{\theta} \mathbf{e}_{2} e^{i\phi}
= C_{\theta} \hat{\boldsymbol{\phi}}.$$
(1.10)

The $\hat{\phi}$ partials are

$$\frac{\partial \hat{\boldsymbol{\phi}}}{\partial \theta} = \frac{\partial}{\partial \theta} \mathbf{e}_2 e^{i\phi}$$

$$= 0.$$
(1.11)

$$\frac{\partial \hat{\boldsymbol{\phi}}}{\partial \boldsymbol{\phi}} = \frac{\partial}{\partial \boldsymbol{\phi}} \mathbf{e}_{2} e^{i\boldsymbol{\phi}}
= -\mathbf{e}_{1} e^{i\boldsymbol{\phi}}
= -\hat{\mathbf{r}} \left\langle \hat{\mathbf{r}} \mathbf{e}_{1} e^{i\boldsymbol{\phi}} \right\rangle - \hat{\boldsymbol{\theta}} \left\langle \hat{\boldsymbol{\theta}} \mathbf{e}_{1} e^{i\boldsymbol{\phi}} \right\rangle - \hat{\boldsymbol{\phi}} \left\langle \hat{\boldsymbol{\phi}} \mathbf{e}_{1} e^{i\boldsymbol{\phi}} \right\rangle
= -\hat{\mathbf{r}} \left\langle \left(\mathbf{e}_{1} e^{i\boldsymbol{\phi}} S_{\theta} + \mathbf{e}_{3} C_{\theta} \right) \mathbf{e}_{1} e^{i\boldsymbol{\phi}} \right\rangle - \hat{\boldsymbol{\theta}} \left\langle \left(C_{\theta} \mathbf{e}_{1} e^{i\boldsymbol{\phi}} - S_{\theta} \mathbf{e}_{3} \right) \mathbf{e}_{1} e^{i\boldsymbol{\phi}} \right\rangle
= -\hat{\mathbf{r}} \left\langle e^{-i\boldsymbol{\phi}} S_{\theta} e^{i\boldsymbol{\phi}} \right\rangle - \hat{\boldsymbol{\theta}} \left\langle C_{\theta} e^{-i\boldsymbol{\phi}} e^{i\boldsymbol{\phi}} \right\rangle
= -\hat{\mathbf{r}} S_{\theta} - \hat{\boldsymbol{\theta}} C_{\theta}.$$
(1.12)

The $\hat{\mathbf{r}}$ partials are were computed as a side effect of evaluating \mathbf{x}_{θ} , and \mathbf{x}_{ϕ} , and are

$$\frac{\partial \hat{\mathbf{r}}}{\partial \theta} = \hat{\boldsymbol{\theta}},\tag{1.13}$$

$$\frac{\partial \hat{\mathbf{r}}}{\partial \phi} = S_{\theta} \hat{\boldsymbol{\phi}}. \tag{1.14}$$

In summary

$$\partial_{\theta}\hat{\mathbf{r}} = \hat{\boldsymbol{\theta}}
\partial_{\phi}\hat{\mathbf{r}} = S_{\theta}\hat{\boldsymbol{\phi}}
\partial_{\theta}\hat{\boldsymbol{\theta}} = -\hat{\mathbf{r}}
\partial_{\phi}\hat{\boldsymbol{\theta}} = C_{\theta}\hat{\boldsymbol{\phi}}
\partial_{\theta}\hat{\boldsymbol{\phi}} = 0
\partial_{\phi}\hat{\boldsymbol{\phi}} = -\hat{\mathbf{r}}S_{\theta} - \hat{\boldsymbol{\theta}}C_{\theta}.$$
(1.15)

Divergence and curl. The divergence and curl can be computed from the vector product of the spherical coordinate gradient and the spherical representation of a vector. That is

$$\nabla \mathbf{A} = \nabla \cdot \mathbf{A} + \nabla \wedge \mathbf{A} = \nabla \cdot \mathbf{A} + I \nabla \times \mathbf{A}. \tag{1.16}$$

That gradient vector product is

$$\nabla \mathbf{A} = \left(\hat{\mathbf{r}}\partial_{r} + \frac{\hat{\boldsymbol{\theta}}}{r}\partial_{\theta} + \frac{\hat{\boldsymbol{\phi}}}{rS_{\theta}}\partial_{\phi}\right) \left(\hat{\mathbf{r}}A_{r} + \hat{\boldsymbol{\theta}}A_{\theta} + \hat{\boldsymbol{\phi}}A_{\phi}\right)$$

$$= \hat{\mathbf{r}}\partial_{r} \left(\hat{\mathbf{r}}A_{r} + \hat{\boldsymbol{\theta}}A_{\theta} + \hat{\boldsymbol{\phi}}A_{\phi}\right)$$

$$+ \frac{\hat{\boldsymbol{\theta}}}{r}\partial_{\theta} \left(\hat{\mathbf{r}}A_{r} + \hat{\boldsymbol{\theta}}A_{\theta} + \hat{\boldsymbol{\phi}}A_{\phi}\right)$$

$$+ \frac{\hat{\boldsymbol{\phi}}}{rS_{\theta}}\partial_{\hat{\boldsymbol{\phi}}} \left(\hat{\mathbf{r}}A_{r} + \hat{\boldsymbol{\theta}}A_{\theta} + \hat{\boldsymbol{\phi}}A_{\phi}\right)$$

$$= \left(\partial_{r}A_{r} + \hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\partial_{r}A_{\theta} + \hat{\mathbf{r}}\hat{\boldsymbol{\phi}}\partial_{r}A_{\phi}\right)$$

$$+ \frac{1}{r} \left(\hat{\boldsymbol{\theta}}(\partial_{\theta}\hat{\mathbf{r}})A_{r} + \hat{\boldsymbol{\theta}}(\partial_{\theta}\hat{\boldsymbol{\theta}})A_{\theta} + \hat{\boldsymbol{\theta}}(\partial_{\theta}\hat{\boldsymbol{\phi}})A_{\phi} + \hat{\boldsymbol{\theta}}\hat{\mathbf{r}}\partial_{\theta}A_{r} + \partial_{\theta}A_{\theta} + \hat{\boldsymbol{\theta}}\hat{\boldsymbol{\phi}}\partial_{\theta}A_{\phi}\right)$$

$$+ \frac{1}{rS_{\theta}} \left(\hat{\boldsymbol{\phi}}(\partial_{\phi}\hat{\mathbf{r}})A_{r} + \hat{\boldsymbol{\phi}}(\partial_{\phi}\hat{\boldsymbol{\theta}})A_{\theta} + \hat{\boldsymbol{\phi}}(\partial_{\phi}\hat{\boldsymbol{\phi}})A_{\phi} + \hat{\boldsymbol{\phi}}\hat{\mathbf{r}}\partial_{\phi}A_{r} + \hat{\boldsymbol{\phi}}\hat{\boldsymbol{\theta}}\partial_{\phi}A_{\theta} + \partial_{\phi}A_{\phi}\right)$$

$$= \left(\partial_{r}A_{r} + \hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\partial_{r}A_{\theta} + \hat{\mathbf{r}}\hat{\boldsymbol{\phi}}\partial_{r}A_{\phi}\right)$$

$$+ \frac{1}{r} \left(\hat{\boldsymbol{\theta}}(\hat{\boldsymbol{\theta}})A_{r} + \hat{\boldsymbol{\theta}}(-\hat{\mathbf{r}})A_{\theta} + \hat{\boldsymbol{\theta}}(0)A_{\phi} + \hat{\boldsymbol{\theta}}\hat{\mathbf{r}}\partial_{\theta}A_{r} + \partial_{\theta}A_{\theta} + \hat{\boldsymbol{\theta}}\hat{\boldsymbol{\phi}}\partial_{\theta}A_{\phi}\right)$$

$$+ \frac{1}{rS_{\theta}} \left(\hat{\boldsymbol{\phi}}(S_{\theta}\hat{\boldsymbol{\phi}})A_{r} + \hat{\boldsymbol{\phi}}(C_{\theta}\hat{\boldsymbol{\phi}})A_{\theta} - \hat{\boldsymbol{\phi}}(\hat{\mathbf{r}}S_{\theta} + \hat{\boldsymbol{\theta}}C_{\theta})A_{\phi} + \hat{\boldsymbol{\phi}}\hat{\mathbf{r}}\partial_{\phi}A_{r} + \hat{\boldsymbol{\phi}}\hat{\boldsymbol{\theta}}\partial_{\phi}A_{\theta} + \partial_{\phi}A_{\phi}\right).$$

The scalar component of this is the divergence

$$\nabla \cdot \mathbf{A} = \partial_{r} A_{r} + \frac{A_{r}}{r} + \frac{1}{r} \partial_{\theta} A_{\theta} + \frac{1}{rS_{\theta}} \left(S_{\theta} A_{r} + C_{\theta} A_{\theta} + \partial_{\phi} A_{\phi} \right)$$

$$= \partial_{r} A_{r} + 2 \frac{A_{r}}{r} + \frac{1}{r} \partial_{\theta} A_{\theta} + \frac{1}{rS_{\theta}} C_{\theta} A_{\theta} + \frac{1}{rS_{\theta}} \partial_{\phi} A_{\phi}$$

$$= \partial_{r} A_{r} + 2 \frac{A_{r}}{r} + \frac{1}{r} \partial_{\theta} A_{\theta} + \frac{1}{rS_{\theta}} C_{\theta} A_{\theta} + \frac{1}{rS_{\theta}} \partial_{\phi} A_{\phi},$$

$$(1.18)$$

which can be factored as

$$\nabla \cdot \mathbf{A} = \frac{1}{r^2} \partial_r (r^2 A_r) + \frac{1}{r S_\theta} \partial_\theta (S_\theta A_\theta) + \frac{1}{r S_\theta} \partial_\phi A_\phi.$$
 (1.19)

The bivector grade of ∇A is the bivector curl

$$\nabla \wedge \mathbf{A} = \left(\hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\partial_{r}A_{\theta} + \hat{\mathbf{r}}\hat{\boldsymbol{\phi}}\partial_{r}A_{\phi}\right) + \frac{1}{r}\left(\hat{\boldsymbol{\theta}}(-\hat{\mathbf{r}})A_{\theta} + \hat{\boldsymbol{\theta}}\hat{\mathbf{r}}\partial_{\theta}A_{r} + \hat{\boldsymbol{\theta}}\hat{\boldsymbol{\phi}}\partial_{\theta}A_{\phi}\right)$$

$$+ \frac{1}{rS_{\theta}}\left(-\hat{\boldsymbol{\phi}}(\hat{\mathbf{r}}S_{\theta} + \hat{\boldsymbol{\theta}}C_{\theta})A_{\phi} + \hat{\boldsymbol{\phi}}\hat{\mathbf{r}}\partial_{\phi}A_{r} + \hat{\boldsymbol{\phi}}\hat{\boldsymbol{\theta}}\partial_{\phi}A_{\theta}\right)$$

$$= \left(\hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\partial_{r}A_{\theta} - \hat{\boldsymbol{\phi}}\hat{\mathbf{r}}\partial_{r}A_{\phi}\right) + \frac{1}{r}\left(\hat{\mathbf{r}}\hat{\boldsymbol{\theta}}A_{\theta} - \hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\partial_{\theta}A_{r} + \hat{\boldsymbol{\theta}}\hat{\boldsymbol{\phi}}\partial_{\theta}A_{\phi}\right)$$

$$+ \frac{1}{rS_{\theta}}\left(-\hat{\boldsymbol{\phi}}\hat{\mathbf{r}}S_{\theta}A_{\phi} + \hat{\boldsymbol{\theta}}\hat{\boldsymbol{\phi}}C_{\theta}A_{\phi} + \hat{\boldsymbol{\phi}}\hat{\mathbf{r}}\partial_{\phi}A_{r} - \hat{\boldsymbol{\theta}}\hat{\boldsymbol{\phi}}\partial_{\phi}A_{\theta}\right)$$

$$= \hat{\boldsymbol{\theta}}\hat{\boldsymbol{\phi}}\left(\frac{1}{rS_{\theta}}C_{\theta}A_{\phi} + \frac{1}{r}\partial_{\theta}A_{\phi} - \frac{1}{rS_{\theta}}\partial_{\phi}A_{\theta}\right)$$

$$+ \hat{\boldsymbol{\phi}}\hat{\mathbf{r}}\left(-\partial_{r}A_{\phi} + \frac{1}{r}\partial_{\theta}A_{\phi} - \frac{1}{rS_{\theta}}\partial_{\phi}A_{r}\right) + \hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\left(\partial_{r}A_{\theta} + \frac{1}{r}A_{\theta} - \frac{1}{r}\partial_{\theta}A_{r}\right)$$

$$= I\hat{\mathbf{r}}\left(\frac{1}{rS_{\theta}}\partial_{\theta}(S_{\theta}A_{\phi}) - \frac{1}{rS_{\theta}}\partial_{\phi}A_{\theta}\right) + I\hat{\boldsymbol{\theta}}\left(\frac{1}{rS_{\theta}}\partial_{\phi}A_{r} - \frac{1}{r}\partial_{r}(rA_{\phi})\right) + I\hat{\boldsymbol{\phi}}\left(\frac{1}{r}\partial_{r}(rA_{\theta}) - \frac{1}{r}\partial_{\theta}A_{r}\right)$$

$$(1.20)$$

This gives

$$\nabla \times \mathbf{A} = \hat{\mathbf{r}} \left(\frac{1}{rS_{\theta}} \partial_{\theta} (S_{\theta} A_{\phi}) - \frac{1}{rS_{\theta}} \partial_{\phi} A_{\theta} \right) + \hat{\boldsymbol{\theta}} \left(\frac{1}{rS_{\theta}} \partial_{\phi} A_r - \frac{1}{r} \partial_r (rA_{\phi}) \right) + \hat{\boldsymbol{\phi}} \left(\frac{1}{r} \partial_r (rA_{\theta}) - \frac{1}{r} \partial_{\theta} A_r \right).$$

$$(1.21)$$

This and the divergence result above both check against the back cover of [1].

Laplacian Using the divergence and curl it's possible to compute the Laplacian from those, but we saw in cylindrical coordinates that it was much harder to do it that way than to do it directly.

$$\nabla^{2}\psi = \left(\hat{\mathbf{r}}\partial_{r} + \frac{\hat{\boldsymbol{\theta}}}{r}\partial_{\theta} + \frac{\hat{\boldsymbol{\phi}}}{rS_{\theta}}\partial_{\phi}\right) \left(\hat{\mathbf{r}}\partial_{r}\psi + \frac{\hat{\boldsymbol{\theta}}}{r}\partial_{\theta}\psi + \frac{\hat{\boldsymbol{\phi}}}{rS_{\theta}}\partial_{\phi}\psi\right)$$

$$= \partial_{rr}\psi + \hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\partial_{r} \left(\frac{1}{r}\partial_{\theta}\psi\right) + \hat{\mathbf{r}}\hat{\boldsymbol{\phi}}\frac{1}{S_{\theta}}\partial_{r} \left(\frac{1}{r}\partial_{\phi}\psi\right)$$

$$+ \frac{\hat{\boldsymbol{\theta}}}{r}\partial_{\theta} \left(\hat{\mathbf{r}}\partial_{r}\psi\right) + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta} \left(\hat{\boldsymbol{\theta}}\partial_{\theta}\psi\right) + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta} \left(\frac{\hat{\boldsymbol{\phi}}}{S_{\theta}}\partial_{\phi}\psi\right)$$

$$+ \frac{\hat{\boldsymbol{\phi}}}{rS_{\theta}}\partial_{\phi} \left(\hat{\mathbf{r}}\partial_{r}\psi\right) + \frac{\hat{\boldsymbol{\phi}}}{r^{2}S_{\theta}}\partial_{\phi} \left(\hat{\boldsymbol{\theta}}\partial_{\theta}\psi\right) + \frac{\hat{\boldsymbol{\phi}}}{r^{2}S_{\theta}^{2}}\partial_{\phi} \left(\hat{\boldsymbol{\phi}}\partial_{\phi}\psi\right)$$

$$= \partial_{rr}\psi + \hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\partial_{r} \left(\frac{1}{r}\partial_{\theta}\psi\right) + \hat{\mathbf{r}}\hat{\boldsymbol{\phi}}\frac{1}{S_{\theta}}\partial_{r} \left(\frac{1}{r}\partial_{\phi}\psi\right)$$

$$+ \frac{\hat{\boldsymbol{\theta}}}{r}\partial_{\theta} \left(\partial_{r}\psi\right) + \frac{1}{r^{2}}\partial_{\theta\theta}\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta} \left(\frac{1}{S_{\theta}}\partial_{\phi}\psi\right)$$

$$+ \frac{\hat{\boldsymbol{\phi}}}{r}\partial_{\theta} \left(\partial_{r}\psi\right) + \frac{\hat{\boldsymbol{\phi}}}{r^{2}}\partial_{\theta}\partial_{\phi}\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\partial_{\phi}\psi$$

$$+ \frac{\hat{\boldsymbol{\theta}}}{r}(\partial_{\theta}\hat{\mathbf{r}})\partial_{r}\psi + \frac{\hat{\boldsymbol{\phi}}}{r^{2}}(\partial_{\theta}\hat{\boldsymbol{\theta}})\partial_{\theta}\psi + \frac{\hat{\boldsymbol{\phi}}}{r^{2}}\partial_{\theta}\partial_{\phi}\psi$$

$$+ \frac{\hat{\boldsymbol{\phi}}}{r}(\partial_{\theta}\hat{\mathbf{r}})\partial_{r}\psi + \frac{\hat{\boldsymbol{\phi}}}{r^{2}}(\partial_{\theta}\hat{\boldsymbol{\theta}})\partial_{\theta}\psi + \frac{\hat{\boldsymbol{\phi}}}{r^{2}}\partial_{\theta}\partial_{\phi}\psi$$

$$= \partial_{rr}\psi + \hat{\mathbf{r}}\hat{\boldsymbol{\theta}}\partial_{r} \left(\frac{1}{r}\partial_{\theta}\psi\right) + \hat{\mathbf{r}}\hat{\boldsymbol{\phi}}\frac{1}{S_{\theta}}\partial_{r} \left(\frac{1}{r}\partial_{\phi}\psi\right)$$

$$+ \frac{\hat{\boldsymbol{\theta}}\hat{\boldsymbol{\tau}}}{r}\partial_{\theta} \left(\partial_{r}\psi\right) + \frac{1}{r^{2}}\partial_{\theta\theta}\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta} \left(\frac{1}{r}\partial_{\phi}\psi\right)$$

$$+ \frac{\hat{\boldsymbol{\phi}}\hat{\boldsymbol{\tau}}}{r}\partial_{\theta} \left(\partial_{r}\psi\right) + \frac{1}{r^{2}}\partial_{\theta\theta}\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta} \left(\frac{1}{r}\partial_{\phi}\psi\right)$$

$$+ \frac{\hat{\boldsymbol{\theta}}\hat{\boldsymbol{\tau}}}{r}\partial_{\theta} \left(\partial_{r}\psi\right) + \frac{1}{r^{2}}\partial_{\theta\theta}\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta} \left(\frac{1}{S_{\theta}}\partial_{\phi}\psi\right)$$

$$+ \frac{\hat{\boldsymbol{\theta}}\hat{\boldsymbol{\tau}}}{r}\partial_{\theta} \left(\partial_{r}\psi\right) + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi$$

$$+ \frac{\hat{\boldsymbol{\theta}}\hat{\boldsymbol{\tau}}}{r}\partial_{\theta} \left(\partial_{r}\psi\right) + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi$$

$$+ \frac{\hat{\boldsymbol{\theta}}\hat{\boldsymbol{\tau}}}{r}\partial_{\theta} \left(\partial_{r}\psi\right) + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi$$

$$+ \frac{\hat{\boldsymbol{\theta}}\hat{\boldsymbol{\tau}}}{r}\partial_{\theta} \left(\partial_{r}\psi\right) + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi$$

$$+ \frac{\hat{\boldsymbol{\theta}}\hat{\boldsymbol{\tau}}}{r}\partial_{\theta} \left(\partial_{r}\psi\right) + \frac{\hat{\boldsymbol{\theta}}\hat{\boldsymbol{\tau}}\partial_{\theta}\psi\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta}\psi\psi + \frac{\hat{\boldsymbol{\theta}}}{r^{2}}\partial_{\theta$$

All the bivector factors are expected to cancel out, but this should be checked. Those with an $\hat{r}\hat{\theta}$ factor are

$$\partial_r \left(\frac{1}{r} \partial_\theta \psi \right) - \frac{1}{r} \partial_{\theta r} \psi + \frac{1}{r^2} \partial_\theta \psi = -\frac{1}{r^2} \partial_\theta \psi + \frac{1}{r} \partial_{r\theta} \psi - \frac{1}{r} \partial_{\theta r} \psi + \frac{1}{r^2} \partial_\theta \psi$$

$$= 0. \tag{1.23}$$

and those with a $\hat{\theta}\hat{\phi}$ factor are

$$\frac{1}{r^2}\partial_{\theta}\left(\frac{1}{S_{\theta}}\partial_{\phi}\psi\right) - \frac{1}{r^2S_{\theta}}\partial_{\phi\theta}\psi + \frac{1}{r^2S_{\theta}^2}C_{\theta}\partial_{\phi}\psi = -\frac{1}{r^2}\frac{C_{\theta}}{S_{\theta}^2}\partial_{\phi}\psi + \frac{1}{r^2S_{\theta}}\partial_{\theta\phi}\psi - \frac{1}{r^2S_{\theta}}\partial_{\phi\theta}\psi + \frac{1}{r^2S_{\theta}^2}C_{\theta}\partial_{\phi}\psi = 0,$$

$$= 0,$$

$$(1.24)$$

and those with a $\hat{\phi}\hat{\mathbf{r}}$ factor are

$$-\frac{1}{S_{\theta}}\partial_{r}\left(\frac{1}{r}\partial_{\phi}\psi\right) + \frac{1}{rS_{\theta}}\partial_{\phi r}\psi - \frac{1}{r^{2}S_{\theta}^{2}}S_{\theta}\partial_{\phi}\psi = \frac{1}{S_{\theta}}\frac{1}{r^{2}}\partial_{\phi}\psi - \frac{1}{rS_{\theta}}\partial_{r\phi}\psi + \frac{1}{rS_{\theta}}\partial_{\phi r}\psi - \frac{1}{r^{2}S_{\theta}}\partial_{\phi}\psi$$

$$= 0.$$

$$(1.25)$$

This leaves

$$\nabla^2 \psi = \partial_{rr} \psi + \frac{2}{r} \partial_r \psi + \frac{1}{r^2} \partial_{\theta\theta} \psi + \frac{1}{r^2 S_\theta} C_\theta \partial_\theta \psi + \frac{1}{r^2 S_\theta^2} \partial_{\phi\phi} \psi. \tag{1.26}$$

This factors nicely as

$$\nabla^2 \psi = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2 \sin \theta^2} \frac{\partial^2 \psi}{\partial \phi^2}, \tag{1.27}$$

which checks against the back cover of Jackson. Here it has been demonstrated explicitly that this operator expression is valid for multivector fields ψ as well as scalar fields ψ .

Bibliography

- [1] JD Jackson. Classical Electrodynamics. John Wiley and Sons, 2nd edition, 1975. 1
- [2] A. Macdonald. *Vector and Geometric Calculus*. CreateSpace Independent Publishing Platform, 2012. 1