

Momentum of scalar field

Way back in lecture 8, it was claimed that

$$P^k = \int d^3x \hat{\pi} \partial^k \hat{\phi} = \int \frac{d^3p}{(2\pi)^3} p^k a_{\mathbf{p}}^\dagger a_{\mathbf{p}}. \quad (1.1)$$

If I compute this, I get a normal ordered variation of this operator, but also get some time dependent terms. Here's the computation (dropping hats)

$$\begin{aligned} P^k &= \int d^3x \hat{\pi} \partial^k \phi \\ &= \int d^3x \partial_0 \phi \partial^k \phi \\ &= \int d^3x \frac{d^3p d^3q}{(2\pi)^6} \frac{1}{\sqrt{2\omega_p 2\omega_q}} \partial_0 \left(a_{\mathbf{p}} e^{-ip \cdot x} + a_{\mathbf{p}}^\dagger e^{ip \cdot x} \right) \partial^k \left(a_{\mathbf{q}} e^{-iq \cdot x} + a_{\mathbf{q}}^\dagger e^{iq \cdot x} \right). \end{aligned} \quad (1.2)$$

The exponential derivatives are

$$\begin{aligned} \partial_0 e^{\pm ip \cdot x} &= \partial_0 e^{\pm i p_\mu x^\mu} \\ &= \pm i p_0 \partial_0 e^{\pm ip \cdot x}, \end{aligned} \quad (1.3)$$

and

$$\begin{aligned} \partial^k e^{\pm ip \cdot x} &= \partial^k e^{\pm i p^\mu x_\mu} \\ &= \pm i p^k e^{\pm ip \cdot x}, \end{aligned} \quad (1.4)$$

so

$$\begin{aligned} P^k &= - \int d^3x \frac{d^3p d^3q}{(2\pi)^6} \frac{1}{\sqrt{2\omega_p 2\omega_q}} p_0 q^k \left(-a_{\mathbf{p}} e^{-ip \cdot x} + a_{\mathbf{p}}^\dagger e^{ip \cdot x} \right) \left(-a_{\mathbf{q}} e^{-iq \cdot x} + a_{\mathbf{q}}^\dagger e^{iq \cdot x} \right) \\ &= - \frac{1}{2} \int d^3x \frac{d^3p d^3q}{(2\pi)^6} \sqrt{\frac{\omega_p}{\omega_q}} q^k \left(a_{\mathbf{p}} a_{\mathbf{q}} e^{-i(p+q) \cdot x} + a_{\mathbf{p}}^\dagger a_{\mathbf{q}}^\dagger e^{i(p+q) \cdot x} - a_{\mathbf{p}} a_{\mathbf{q}}^\dagger e^{i(q-p) \cdot x} - a_{\mathbf{p}}^\dagger a_{\mathbf{q}} e^{i(p-q) \cdot x} \right) \\ &= \frac{1}{2} \int \frac{d^3p d^3q}{(2\pi)^3} \sqrt{\frac{\omega_p}{\omega_q}} q^k \left(-a_{\mathbf{p}} a_{\mathbf{q}} e^{-i(\omega_p + \omega_q)t} \delta^3(\mathbf{p} + \mathbf{q}) - a_{\mathbf{p}}^\dagger a_{\mathbf{q}}^\dagger e^{i(\omega_p + \omega_q)t} \delta^3(-\mathbf{p} - \mathbf{q}) \right. \\ &\quad \left. + a_{\mathbf{p}} a_{\mathbf{q}}^\dagger e^{i(\omega_q - \omega_p)t} \delta^3(\mathbf{p} - \mathbf{q}) + a_{\mathbf{p}}^\dagger a_{\mathbf{q}} e^{i(\omega_p - \omega_q)t} \delta^3(\mathbf{q} - \mathbf{p}) \right) \\ &= \frac{1}{2} \int \frac{d^3p}{(2\pi)^3} p^k \left(a_{\mathbf{p}}^\dagger a_{\mathbf{p}} + a_{\mathbf{p}} a_{\mathbf{p}}^\dagger - a_{\mathbf{p}} a_{-\mathbf{p}} e^{-2i\omega_p t} - a_{\mathbf{p}}^\dagger a_{-\mathbf{p}}^\dagger e^{2i\omega_p t} \right). \end{aligned} \quad (1.5)$$

What is the rationale for ignoring those time dependent terms? Does normal ordering also implicitly drop any non-paired creation/annihilation operators? If so, why?