Matrix form for discrete time Fourier transform

1.1 Transform pair

In [2] a verification of the discrete Fourier transform pairs was performed. A much different looking discrete Fourier transform pair is given in [1] §A.4. This transform pair samples the points at what are called the Nykvist time instants given by

$$t_k = \frac{Tk}{2N+1}, \qquad k \in [-N, \cdots N]$$
(1.1)

Note that the endpoints of these sampling points are not $\pm T/2$, but are instead at

$$\pm \frac{T}{2} \frac{1}{1 + 1/N'} \tag{1.2}$$

which are slightly within the interior of the [-T/2, T/2] range of interest. The reason for this slightly odd seeming selection of sampling times becomes clear if one calculate the inversion relations.

Given a periodic ($\omega_0 T = 2\pi$) bandwidth limited signal evaluated only at the Nykvist times t_k ,

$$x(t_k) = \sum_{n=-N}^{N} X_n e^{jn\omega_0 t_k},$$
(1.3)

assume that an inversion relation can be found. To find X_n evaluate the sum

$$\sum_{k=-N}^{N} x(t_k) e^{-jm\omega_0 t_k} = \sum_{k=-N}^{N} \left(\sum_{n=-N}^{N} X_n e^{jn\omega_0 t_k} \right) e^{-jm\omega_0 t_k}$$

$$= \sum_{n=-N}^{N} X_n \sum_{k=-N}^{N} e^{j(n-m)\omega_0 t_k}$$
(1.4)

This interior sum has the value 2N+1 when n=m. For $n\neq m$, and $a=e^{j(n-m)\frac{2\pi}{2N+1}}$, this is

$$\sum_{k=-N}^{N} e^{j(n-m)\omega_0 t_k} = \sum_{k=-N}^{N} e^{j(n-m)\omega_0 \frac{Tk}{2N+1}}$$

$$= \sum_{k=-N}^{N} a^k$$

$$= a^{-N} \sum_{k=-N}^{N} a^{k+N}$$

$$= a^{-N} \sum_{r=0}^{N} a^r$$

$$= a^{-N} \frac{a^{2N+1} - 1}{a-1}.$$
(1.5)

Since $a^{2N+1} = e^{2\pi j(n-m)} = 1$, this sum is zero when $n \neq m$. This means that

$$\sum_{k=-N}^{N} e^{j(n-m)\omega_0 t_k} = (2N+1)\delta_{n,m},\tag{1.6}$$

which provides the desired Fourier inversion relation

$$X_m = \frac{1}{2N+1} \sum_{k=-N}^{N} x(t_k) e^{-jm\omega_0 t_k}.$$
 (1.7)

1.2 Matrix form

The discrete time Fourier transform has been seen to have the form

$$x_k = \sum_{n=-N}^{N} X_n e^{2\pi j nk/(2N+1)}$$
 (1.8a)

$$X_n = \frac{1}{2N+1} \sum_{k=-N}^{N} x_k e^{-2\pi j nk/(2N+1)}.$$
 (1.8b)

A matrix representation of this form is desired. Let

$$\mathbf{x} = \begin{bmatrix} x_{-N} \\ \vdots \\ x_0 \\ \vdots \\ x_N \end{bmatrix}$$
 (1.9a)

$$\mathbf{X} = \begin{bmatrix} X_{-N} \\ \vdots \\ X_0 \\ \vdots \\ X_N \end{bmatrix}$$
 (1.9b)

Equation (1.8a) written out in full is

$$x_{k} = X_{-N}e^{-2\pi jNk/(2N+1)}$$

$$+ X_{1-N}e^{-2\pi j(N-1)k/(2N+1)}$$

$$+ \cdots$$

$$+ X_{0}$$

$$+ \cdots$$

$$+ X_{N-1}e^{2\pi j(N-1)k/(2N+1)}$$

$$+ X_{N}e^{2\pi jNk/(2N+1)}$$

$$(1.10)$$

With $\alpha = e^{2\pi j/(2N+1)}$ the matrix form is

$$\mathbf{x} = \begin{bmatrix} \alpha^{NN} & \alpha^{(N-1)N} & \cdots & 1 & \cdots & \alpha^{-(N-1)N} & \alpha^{-NN} \\ \alpha^{N(N-1)} & \alpha^{(N-1)(N-1)} & \cdots & 1 & \cdots & \alpha^{-(N-1)(N-1)} & \alpha^{-N(N-1)} \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ 1 & 1 & 1 & 1 & 1 & 1 & 1 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ \alpha^{-N(N-1)} & \alpha^{-(N-1)(N-1)} & \cdots & 1 & \cdots & \alpha^{N-1(N-1)} & \alpha^{N(N-1)} \\ \alpha^{-NN} & \alpha^{-NN} & \cdots & 1 & \cdots & \alpha^{(N-1)N} & \alpha^{NN} \end{bmatrix} \mathbf{X}$$
(1.11)

Similarly, from eq. (1.8b), the inverse relation expands out to

$$(2N+1)X_{n} = x_{-N}e^{2\pi jnN/(2N+1)}$$

$$+ x_{1-N}e^{2\pi jn(N-1)/(2N+1)}$$

$$\cdot \cdot \cdot$$

$$+ x_{0}$$

$$\cdot \cdot \cdot$$

$$+ x_{N-1}e^{-2\pi jn(N-1)/(2N+1)}$$

$$+ x_{N}e^{-2\pi jnN/(2N+1)},$$

$$(1.12)$$

with a matrix form of

$$(2N+1)\mathbf{X} = \begin{bmatrix} \alpha^{-NN} & \alpha^{-N(N-1)} & \cdots & 1 & \cdots & \alpha^{N(N-1)} & \alpha^{NN} \\ \alpha^{-(N-1)N} & \alpha^{-(N-1)(N-1)} & \cdots & 1 & \cdots & \alpha^{(N-1)(N-1)} & \alpha^{(N-1)N} \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ 1 & 1 & 1 & 1 & 1 & 1 & 1 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ \alpha^{(N-1)N} & \alpha^{(N-1)(N-1)} & \cdots & 1 & \cdots & \alpha^{-(N-1)(N-1)} & \alpha^{-(N-1)N} \\ \alpha^{NN} & \alpha^{N(N-1)} & \cdots & 1 & \cdots & \alpha^{-N(N-1)} & \alpha^{-NN} \end{bmatrix}$$
 (1.13)

Letting

$$\mathbf{F} = \begin{bmatrix} \alpha^{NN} & \alpha^{(N-1)N} & \cdots & 1 & \cdots & \alpha^{-(N-1)N} & \alpha^{-NN} \\ \alpha^{N(N-1)} & \alpha^{(N-1)(N-1)} & \cdots & 1 & \cdots & \alpha^{-(N-1)(N-1)} & \alpha^{-N(N-1)} \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ 1 & 1 & 1 & 1 & 1 & 1 & 1 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ \alpha^{-N(N-1)} & \alpha^{-(N-1)(N-1)} & \cdots & 1 & \cdots & \alpha^{N-1(N-1)} & \alpha^{N(N-1)} \\ \alpha^{-NN} & \alpha^{-NN} & \cdots & 1 & \cdots & \alpha^{(N-1)N} & \alpha^{NN} \end{bmatrix},$$
(1.14)

the discrete transform pair has the following compactly matrix representation

$$\mathbf{x} = \mathbf{FX} \tag{1.15a}$$

$$\mathbf{X} = \frac{1}{2N+1}\bar{\mathbf{F}}\mathbf{x},\tag{1.15b}$$

where $\bar{\mathbf{F}}$ is the complex conjugate of \mathbf{F} .

Bibliography

- [1] Franco Giannini and Giorgio Leuzzi. *Nonlinear Microwave Circuit Design*. Wiley Online Library, 2004. 1.1
- [2] Peeter Joot. *Condensed matter physics.*, chapter Discrete Fourier transform. 2013. URL http://peeterjoot.com/archives/math2013/phy487.pdf. [Online; accessed 02-December-2014]. 1.1