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Duality transformation of the far field fields.

We've seen that the far field electric and magnetic fields associated with a magnetic vector potential were

$$\mathbf{E} = -j\omega \operatorname{Proj}_{\mathrm{T}} \mathbf{A},\tag{1.1a}$$

$$\mathbf{H} = \frac{1}{\eta} \hat{\mathbf{k}} \times \mathbf{E}.$$
 (1.1b)

It's worth a quick note that the duality transformation for this, referring to [1] tab.3.2, is

$$\mathbf{H} = -j\omega \operatorname{Proj}_{\mathbf{T}} \mathbf{F}$$
(1.2a)

$$\mathbf{E} = -\eta \hat{\mathbf{k}} \times \mathbf{H}.$$
 (1.2b)

What does **H** look like in terms of **A**, and **E** look like in terms of **H**? The first is

$$\mathbf{H} = -\frac{j\omega}{\eta}\hat{\mathbf{k}} \times \left(\mathbf{A} - \left(\mathbf{A} \cdot \hat{\mathbf{k}}\right)\hat{\mathbf{k}}\right), \qquad (1.3)$$

in which the $\hat{\mathbf{k}}$ crossed terms are killed, leaving

$$\mathbf{H} = -\frac{j\omega}{\eta} \hat{\mathbf{k}} \times \mathbf{A}.$$
 (1.4)

The electric field follows again using a duality transformation, so in terms of the electric vector potential, is

$$\mathbf{E} = j\omega\eta \hat{\mathbf{k}} \times \mathbf{F}.$$
 (1.5)

These show explicitly that neither the electric or magnetic far field have any radial component, matching with intuition for transverse propagation of the fields.

Bibliography

[1] Constantine A Balanis. Antenna theory: analysis and design. John Wiley & Sons, 3rd edition, 2005.