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Canonical bivectors in spacetime algebra.

I've been enjoying XylyXylyX's QED Prerequisites Geometric Algebra: Spacetime YouTube series, which is doing a thorough walk through of [1], filling in missing details. The last episode QED Prerequisites Geometric Algebra 15: Complex Structure, left things with a bit of a cliff hanger, mentioning a "canonical" form for STA bivectors that was intriguing.

The idea is that STA bivectors, like spacetime vectors can be spacelike, timelike, or lightlike (i.e.: positive, negative, or zero square), but can also have a complex signature (squaring to a 0,4-multivector.)

The only context that I knew of that one wanted to square an STA bivector is for the electrodynamic field Lagrangian, which has an F^2 term. In no other context, was the signature of F, the electrodynamic field, of interest that I knew of, so I'd never considered this "Canonical form" representation.

Here are some examples:

$$F = \gamma_{10}, \quad F^{2} = 1$$

$$F = \gamma_{23}, \quad F^{2} = -1$$

$$F = 4\gamma_{10} + \gamma_{13}, \quad F^{2} = 15$$

$$F = \gamma_{10} + \gamma_{13}, \quad F^{2} = 0$$

$$F = \gamma_{10} + 4\gamma_{13}, \quad F^{2} = -15$$

$$F = \gamma_{10} + \gamma_{23}, \quad F^{2} = 2I$$

$$F = \gamma_{10} - 2\gamma_{23}, \quad F^{2} = -3 + 4I.$$
(1.1)

You can see in this table that all the *F*'s that are purely electric, have a positive signature, and all the purely magnetic fields have a negative signature, but when there is a mix, anything goes. The idea behind the canonical representation in the paper is to write

$$F = f e^{l\phi}, \tag{1.2}$$

where f^2 is real and positive, assuming that *F* is not lightlike.

The paper gives a formula for computing *f* and ϕ , but let's do this by example, putting all the F^{2} 's

above into their complex polar form representation, like so

$$F = \gamma_{10}, \quad F^{2} = 1$$

$$F = \gamma_{23}, \quad F^{2} = 1e^{\pi I}$$

$$F = 4\gamma_{10} + \gamma_{13}, \quad F^{2} = 15$$

$$F = \gamma_{10} + \gamma_{13}, \quad F^{2} = 0$$

$$F = \gamma_{10} + 4\gamma_{13}, \quad F^{2} = 15e^{\pi I}$$

$$F = \gamma_{10} + \gamma_{23}, \quad F^{2} = 2e^{(\pi/2)I}$$

$$F = \gamma_{10} - 2\gamma_{23}, \quad F^{2} = 5e^{(\pi - \arctan(4/3))I}$$
(1.3)

Since we can put F^2 in polar form, we can factor out half of that phase angle, so that we are left with a bivector that has a positive square. If we write

$$F^2 = |F^2|e^{2\phi I}, (1.4)$$

we can then form

$$f = F e^{-\phi I}.$$
(1.5)

If we want an equation for ϕ , we can just write

$$2\phi = \operatorname{Arg}(F^2). \tag{1.6}$$

This is a bit better (I think) than the form given in the paper, since it will uniformly rotate F^2 toward the positive region of the real axis, whereas the paper's formula sometimes rotates towards the negative reals, which is a strange seeming polar form to use.

Let's compute *f* for $F = \gamma_{10} - 2\gamma_{23}$, using

$$2\phi = \pi - \arctan(4/3). \tag{1.7}$$

The exponential expands to

$$e^{-\phi I} = \frac{1}{\sqrt{5}} \left(1 - 2I \right). \tag{1.8}$$

Multiplying each of the bivector components by 1 - 2I, we find

$$\gamma_{10} (1 - 2I) = \gamma_{10} - 2\gamma_{100123}$$

= $\gamma_{10} - 2\gamma_{1123}$
= $\gamma_{10} + 2\gamma_{23}$, (1.9)

and

$$-2\gamma_{23} (1 - 2I) = -2\gamma_{23} + 4\gamma_{230123}$$

= $-2\gamma_{23} + 4\gamma_{23}^2\gamma_{01}$ (1.10)
= $-2\gamma_{23} + 4\gamma_{10}$,

leaving

$$f = \sqrt{5}\gamma_{10},\tag{1.11}$$

so the canonical form is

$$F = \gamma_{10} - 2\gamma_{23} = \sqrt{5}\gamma_{10}\frac{1+2I}{\sqrt{5}}.$$
(1.12)

It's interesting here that f, in this case, is a spatial bivector (i.e.: pure electric field), but that clearly isn't always going to be the case, since we can have a case like,

$$F = 4\gamma_{10} + \gamma_{13} = 4\gamma_{10} + \gamma_{20}I, \qquad (1.13)$$

from the table above, that has both electric and magnetic field components, yet is already in the canonical form, with $F^2 = 15$. The canonical f, despite having a positive square, is not necessarily a spatial bivector (as it may have both grades 1,2 in the spatial representation, not just the electric field, spatial grade-1 component.)

Bibliography

[1] Justin Dressel, Konstantin Y Bliokh, and Franco Nori. Spacetime algebra as a powerful tool for electromagnetism. *Physics Reports*, 589:1–71, 2015. 1