
Verifying the GA form for the symmetric and antisymmetric components of the different rate of strain.

We found geometric algebra representations for the symmetric and antisymmetric components for a gradient-vector direct product. In particular, given

$$d\mathbf{v} = d\mathbf{x} \cdot (\nabla \otimes \mathbf{v}) \quad (1.1)$$

we found

$$\begin{aligned} d\mathbf{x} \cdot \mathbf{d} &= \frac{1}{2} d\mathbf{x} \cdot \left(\nabla \otimes \mathbf{v} + (\nabla \otimes \mathbf{v})^\dagger \right) \\ &= \frac{1}{2} (d\mathbf{x} (\nabla \cdot \mathbf{v}) + \langle \nabla d\mathbf{x}\mathbf{v} \rangle_1), \end{aligned} \quad (1.2)$$

and

$$\begin{aligned} d\mathbf{x} \cdot \boldsymbol{\Omega} &= \frac{1}{2} d\mathbf{x} \cdot \left(\nabla \otimes \mathbf{v} - (\nabla \otimes \mathbf{v})^\dagger \right) \\ &= \frac{1}{2} (d\mathbf{x} (\nabla \cdot \mathbf{v}) - \langle d\mathbf{x}\mathbf{v}\nabla \rangle_1). \end{aligned} \quad (1.3)$$

Let's expand each of these in coordinates to verify that these are correct. For the symmetric component, that is

$$\begin{aligned} d\mathbf{x} \cdot \mathbf{d} &= \frac{1}{2} \left(dx_i \partial_j v_j \mathbf{e}_i + \partial_j dx_i v_k \langle \mathbf{e}_j \mathbf{e}_i \mathbf{e}_k \rangle_1 \right) \\ &= \frac{1}{2} dx_i \left(\partial_j v_j \mathbf{e}_i + \partial_j v_k (\delta_{ji} \mathbf{e}_k + (\mathbf{e}_j \wedge \mathbf{e}_i) \cdot \mathbf{e}_k) \right) \\ &= \frac{1}{2} dx_i \left(\partial_j v_j \mathbf{e}_i + \partial_j v_k (\delta_{ji} \mathbf{e}_k + \delta_{ik} \mathbf{e}_j - \delta_{jk} \mathbf{e}_i) \right) \\ &= \frac{1}{2} dx_i \left(\partial_j v_j \mathbf{e}_i + \partial_i v_k \mathbf{e}_k + \partial_j v_i \mathbf{e}_j - \partial_j v_j \mathbf{e}_i \right) \\ &= \frac{1}{2} dx_i \left(\partial_i v_k \mathbf{e}_k + \partial_j v_i \mathbf{e}_j \right) \\ &= dx_i \frac{1}{2} (\partial_i v_j + \partial_j v_i) \mathbf{e}_j. \end{aligned} \quad (1.4)$$

Sure enough, we that the product contains the matrix element of the symmetric component of $\nabla \otimes \mathbf{v}$.

Now let's verify that our GA antisymmetric tensor product representation works out.

$$\begin{aligned}
d\mathbf{x} \cdot \boldsymbol{\Omega} &= \frac{1}{2} \left(dx_i \partial_j v_j \mathbf{e}_i - dx_i \partial_k v_j \langle \mathbf{e}_i \mathbf{e}_j \mathbf{e}_k \rangle_1 \right) \\
&= \frac{1}{2} dx_i (\partial_j v_j \mathbf{e}_i - \partial_k v_j (\delta_{ij} \mathbf{e}_k + \delta_{jk} \mathbf{e}_i - \delta_{ik} \mathbf{e}_j)) \\
&= \frac{1}{2} dx_i (\partial_j v_j \mathbf{e}_i - \partial_k v_i \mathbf{e}_k - \partial_k v_k \mathbf{e}_i + \partial_i v_j \mathbf{e}_j) \\
&= \frac{1}{2} dx_i (\partial_i v_j \mathbf{e}_j - \partial_k v_i \mathbf{e}_k) \\
&= dx_i \frac{1}{2} (\partial_i v_j - \partial_j v_i) \mathbf{e}_j.
\end{aligned} \tag{1.5}$$

As expected, we see that this product contains the matrix element of the antisymmetric component of $\nabla \otimes \mathbf{v}$. We also found previously that $\boldsymbol{\Omega}$ is just a curl, namely

$$\boldsymbol{\Omega} = \frac{1}{2} (\nabla \wedge \mathbf{v}) = \frac{1}{2} (\partial_i v_j) \mathbf{e}_i \wedge \mathbf{e}_j, \tag{1.6}$$

which directly encodes the antisymmetric component of $\nabla \otimes \mathbf{v}$. We can also see that by fully expanding $d\mathbf{x} \cdot \boldsymbol{\Omega}$, which gives

$$\begin{aligned}
d\mathbf{x} \cdot \boldsymbol{\Omega} &= dx_i \frac{1}{2} (\partial_j v_k) \mathbf{e}_i \cdot (\mathbf{e}_j \wedge \mathbf{e}_k) \\
&= dx_i \frac{1}{2} (\partial_j v_k) (\delta_{ij} \mathbf{e}_k - \delta_{ik} \mathbf{e}_j) \\
&= dx_i \frac{1}{2} ((\partial_i v_k) \mathbf{e}_k - (\partial_j v_i) \mathbf{e}_j) \\
&= dx_i \frac{1}{2} (\partial_i v_j - \partial_j v_i) \mathbf{e}_j,
\end{aligned} \tag{1.7}$$

as expected.